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Cygnus and directional DM detection

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Current status of direct detection: SI WIMP-nucleon cross section



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Why is there a neutrino floor? → spectral match between DM and solar neutrinos









Nothing mimics dark matter, including solar neutrinos

Fluorine recoils $[8-50 \text{ keV}_r]$



September 6





O° 300° 240° 180° Galactic longitude, *l*

Subtracting the neutrino background



Orange region is inaccessible without directional information

Fluorine recoil in atmospheric He





What would you use to detect keV-scale nuclear recoils?



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zies,	 Need nm—µm resolution to image tracks No way to transport track topology to a readout Highly scrambled by straggling
blumes	• Tracks much shorter than diffusion scale
n-scale are easily TPC has 90s	• Low target mass





Many alternative approaches to directional detection not taken by Cygnus, but being investigated by others, e.g.

Solid

Crystal defect spectroscopy



+ other techniques in nano-engineered detectors, e.g. graphene, nanotubes, DNA/biomaterial

Liquid (Ar/Xe) Columnar recombination





Focus of Cygnus: gas time projection chamber



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Focus of Cygnus: gas time projection chamber



Angular performance

Everything gets worse at lower energies:

- Decreasing quenching factor, means recoils are harder to detect
- Tracks get shorter \rightarrow harder to measure directions
- Harder to distinguish ER/NRs since tracks are short

→ Energy dependence of directional performance is very important, and needs to be the focus now







Example: angular resolution

Dispersion in measured angles relative to initial recoil direction (=1 rad if there is no correlation and angles are isotropic)



<u>Simulated charge readout comparison</u> To realistically discriminate DM and neutrinos, need angular resolution better than ~30°

Sensitivity

Fluorine-based fill gases have high proton spin $\langle S_p \rangle \rightarrow$ naturally good SD-proton limits





- Good directional sensitivity
- Low threshold O(1) keVr
- ~1000 m³ volume

 10^{4}

Important note: these limits are true discovery limits, i.e. a signal can be <u>confirmed</u> as DM

→ comparison of Cygnus limits with other experiments undersells its potential

Modularity is both necessary, and advantageous

$\mathbf{CYGNUS}\text{-}Nm^3$



 $CYGNUS-10 m^3 module$



CYGNUS-10 Boulby, UK

CYGNUS-HD10 Lead, South Dakota

CYGNO Gran Sasso, Italy

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CYGNUS-Andes Chile/Argentina

in the

CYGNUS-KM Kamioka, Japan

100

17

CYGNUS-OZ Stawell, Aus.

What <u>should</u> the signal look like? → a Gaussian peaking towards Cygnus





$$\exp\left(-\frac{\left(v_{\min} + v_{\text{lab}}(t)\cos\theta\right)^2}{2\sigma_v^2}\right)$$

January 1

Standard prediction based on a few assumptions

• The DM scatters elastically

• The DM velocity distribution is a Gaussian (SHM) $f(\mathbf{v}) \sim \exp\left(-\frac{(\mathbf{v} + \mathbf{v}_{\text{lab}})^2}{2\sigma_{\text{s}}^2}\right)$

• DM-nucleus matrix element does not depend on velocity



Interesting to consider cases where these aren't true, e.g.

Non-standard DM-nucleus kinematics/interactions,

- \rightarrow inelastic DM
- \rightarrow EFT operators with transverse velocity dependence
- → Luminous DM
- → Multi-scatter regime, e.g. superheavy/strongly interacting DM
- Non-Gaussian velocity distributions → The Gaia Sausage/Enceladus → Streams and substructure
- DM fluxes from directions other than Cygnus \rightarrow Supernova-produced DM → Cosmic ray-upscattered DM \rightarrow Boosted DM



Interesting to consider cases where these aren't true, e.g.



Exciting prospects: These are all challenging for nondirectional experiments to probe -> Many opportunities for theoretical studies to determine the full physics potential of a directional TPC network

Directional neutrino measurements



Beyond 10 m³ scale, start to get sizeable rate of solar neutrinos via electron recoils

General physics: Measurement of the Migdal effect → Emission of ~keV electron for very low energy NRs. Important for sub-GeV DM searches, but on shaky ground theoretically as it has never been measured

Could be confirmed directionally, using a small-scale TPC!

- Directional TPC network "Cygnus" is an exciting possibility that appears increasingly plausible
- Primary physics goals are to set limits beyond the neutrino floor, and to provide a convincing confirmation of DM in the event of detection
- Exciting physics case for non-WIMP dark matter, neutrino physics, and general physics measurements, all deserving of further exploration now that the community is converging on a strategy

Summary

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Directional Recoil Detection

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Keywords

nuclear recoils, electron recoils, dark matter, neutrinos, gas time projection chambers, Migdal effect

Abstract

Searches for dark matter-induced recoils have made impressive advances in the last few years. Yet the field is confronted by several outstanding problems. First, the inevitable background of solar neutrinos will soon inhibit the conclusive identification of many dark matter models. Second, and more fundamentally, current experiments have no practical way of confirming a detected signal's galactic origin. The concept of directional detection addresses both of these issues while offering opportunities to study novel dark matter and neutrino-related physics. The concept remains experimentally challenging, but gas time projection chambers are an increasingly attractive option, and when properly configured, would allow directional measurements of both nuclear and electron recoils. In this review, we reassess the required detector performance and survey relevant technologies. Fortuitously, the highly-segmented detectors required to achieve good directionality also enable several fundamental and applied physics measurements. We comment on near-term challenges and how the field could be advanced.

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Detector classes by d

Anisotropic scintillators

- No event-level directions
- Exploits modulation of DM with respect to crystal axes

Columnar recombination

- Event-level 1d directions
- No head/tail
- Direction and energy are not independent

lirectional information Demonstrated R&D Proposed						
Recoil imaging						
— Event-level — —						
Time-integrated recoil imaging	Time-resolved recoil imaging					
 Nuclear emulsions 2d recoil tracks, without head/tail No event times information recorded 	 Gas TPC Head/tail measurable 1d, 2d or 3d Independent energy/ direction measurement 					
 DNA detector 3d recoils without head/tail No event times recorded 	 Crystal defects 3d track topology Head/tail measurable 					

Stawell Underground Physics Laboratory (SUPL)

- + 1.6 km depth
- First underground site in Southern Hemisphere

Gas mixture Pressure [Torr] Density $[kg/m^3]$ W [eV/ion pair] Trans. diffusion $|\mu m/\sqrt{2}$ Long. diffusion $\mu m/\sqrt{c}$ Drift velocity [mm/µs] Mean avalanche gain

and is also used for He:SF_6 mixtures.

	${ m SF}_6$	$\operatorname{He:}{\operatorname{SF}_6}$	$\operatorname{He:}{\operatorname{SF}_6}$
	20	740:20	755:5
	0.16	0.32	0.20
	35.5	38.0	40.0
cm]	116.2	78.6	78.6
cm]	116.2	78.6	78.6
	0.140	0.140	0.140
	9×10^3	9×10^3	9×10^3

TABLE I. Various gas-dependent parameters assumed in the TPC detector simulation. The values are sourced as follows: the W factor for pure SF_6 is from a measurement with alpha particles [310], while the W factors for the He:SF_6 and He:CF_4 mixtures are calculated using Eq.(1) of Ref. [266]. The diffusion values and drift velocity in 20 Torr of pure SF_6 were measured in Ref. [299]. For the He:SF_6 mixtures, no measurements or reliable simulations exist, so we use the 40 Torr pure SF_6 diffusion from Ref. [299] and then assume the electric field can be adjusted to keep the drift velocity constant. The avalanche gain assumed for pure SF_6 has been achieved with THGEMs in Ref. [311] and triple thin GEMs in Ref. [312],

Readout type	Dimensionality	Segmentation $(x \times y)$	Capacitance $[pF]$	$\sigma_{ m noise}~{ m in}~1~{ m \mu s}$	$\mathrm{Threshold}/\sigma_{\mathrm{noise}}$
planar	1d(z)	$10~{ m cm}$ $ imes$ $10~{ m cm}$	3000	$18000 \ e^-$	3.09
wire	2d (yz)	$1 \mathrm{~m~wires}, 2 \mathrm{~mm~pitch}$	0.25	$800~e^-$	4.11
pad	3d (xyz)	$3~\mathrm{mm}$ $ imes$ $3~\mathrm{mm}$	0.25	$375 e^-$	4.77
optical	2d(xyz)	$200~\mu{ m m} imes200~\mu{ m m}$	\mathbf{n}/\mathbf{a}	$2 {\rm photons}$	5.77
strip	3d(xyz)	1 m strips, 200 μ m pitch	500	$2800~e^-$	4.61
pixel	3d(xyz)	$200~\mu\mathrm{m}\times200~\mu\mathrm{m}$	0.012 - 0.200	$42~e^-$	5.77

acceptance, optical transparency, and quantum efficiency.

TABLE II. List of readout-specific parameters that are used in the simulation of each technology we consider here. The capacitance, which determines the noise level, is listed as that for a single detector element. For the optical readout, a yield of 7.2×10^{-6} photons per avalanche electron is used to account for the combined effects of photon yield, geometric optical