

Dark matter and collider phenomenology of an SU(2) triplet and singlet scalar extended standard model

Leon Friedrich¹

lfriedrich@student.unimelb.edu.au

Based on work done in collaboration with N. Bell¹, M. Dolan¹, M. Ramsey-Musolf^{2,3,4}, and R. Volkas¹

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¹ ARC Centre of Excellence for Dark Matter Particle Physics, School of Physics, The University of Melbourne

² Amherst Center for Fundamental Interactions, Department of Physics, University of Massachusetts Amherst

³ Tsung-Dao Lee Institute and School of Physics and Astronomy, Shanghai Jiao Tong University

⁴ Kellogg Radiation Laboratory, California Institute of Technology

- Models with new scalars expand the scalar potential parameter space.
- This may lead to a first-order electroweak phase transition, as required for electroweak baryogenesis.
- If these new scalars are also stable, they will contribute to the dark matter density.
- Is there a relatively simple scalar extension that could explain both the DM density and BAU?

Adding a real singlet scalar S is the simplest possible scalar extension.

- Relic density determined by the Higgs coupling, as it is the only coupling to the SM particles.

A Feynman diagram illustrating the production and decay of a Higgs boson (H). Two dashed lines representing singlet scalars (S) meet at a vertex on the left, with a dashed line labeled H extending to the right. This H line then splits into two solid lines labeled SM (Standard Model particles) on the right.

$$\Rightarrow \langle \sigma_A v \rangle \propto \lambda_{HS}^2$$

- This coupling also determines the direct detection cross section.

A Feynman diagram showing the interaction of a Higgs boson (H) with a nucleus (N) and a singlet scalar (S). A dashed line labeled H enters from the bottom, where it meets a vertex with two dashed lines labeled S extending downwards. From this vertex, a solid line labeled H goes upwards to another vertex, from which two solid lines labeled N (nucleus) extend upwards and outwards.

$$\Rightarrow \sigma_{SI} \propto \lambda_{HS}^2$$

- Severe DM constraints are inconsistent with requirements for a SFO EWPT.

Adding a real SU(2) triplet scalar Σ is the next simplest extension.

- Annihilation rate dominated by gauge couplings.

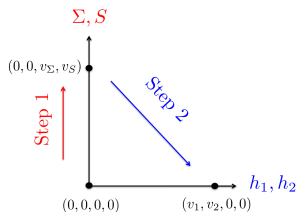
$$\Sigma^0 \quad W^+ \\ \Sigma^0 \quad W^- \quad \Rightarrow \quad \langle \sigma_A v \rangle \propto g^4$$

- This large annihilation rate requires the new scalars to have a very large mass ($m_\Sigma \gtrsim 2$ TeV).
- Similar result for inert doublet DM ($m_{H_2} \gtrsim 500$ GeV).
- When the coupling to the SM Higgs is non-zero, as required for interesting EWPT, these masses need to be even larger.
- Minimal scalar multiplet DM is inconsistent with parameter space required for an interesting EWPT.

Minimal scalar extensions cannot provide a DM candidate while also contributing significantly to the EWPT. Need additional particles.

- Multiple gauge singlet scalars.
- Multiplet $SU(2)$ multiplet scalars, e.g., multiple inert Higgs doublets.
- 2HDM+singlet
- Triplet+Singlet Model

- We examine a model where the SM is extended by a real scalar singlet S and a real scalar triplet $\Sigma \sim (1, 3, 0)$.
- Motivated by a similar model, with two Higgs doublets, that was examined in the context of EWBG (arXiv:1508.05404).



We must impose a \mathbb{Z}_2 symmetry in order to ensure one of the new particles is stable.

- We will charge both the singlet and triplet under a single \mathbb{Z}_2 symmetry.
- Similar models have been examined in arXiv:1311.1077 and arXiv:2009.01262
- The \mathbb{Z}_2 symmetry ensures that there is a stable particle.
- This symmetry also prevents the new scalars from mixing with the SM Higgs.
- However, the singlet and triplet can mix with each other.

- The \mathbb{Z}_2 symmetry permits a $\lambda_{H\Sigma S}H^\dagger\Sigma HS$ term.
- After EWSB this will give rise to mixing between Σ^0 and S
- We rotate to the mass basis

$$\begin{pmatrix} \Sigma^{0'} \\ S' \end{pmatrix} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} \Sigma^0 \\ S \end{pmatrix}.$$

- The heavier scalar will decay into the lighter one and an (off-shell) SM Higgs.
- The lightest scalar will be stable and contribute to the DM density.
- If $m_{\Sigma^{0'}} < m_{S'}$, encounter same issue as in minimal triplet scalar DM, annihilation rate is large \implies need large masses.

Consider $m_{S'} < m_{\Sigma^{0'}}$, S' is the DM. There are three significant processes that could keep the S' in thermal equilibrium.

(a) Equilibrium through weak gauge boson couplings.

A Feynman diagram showing a vertex where a dashed line labeled S' and another dashed line labeled $\Sigma^{0'}$ meet. From this vertex, two wavy lines emerge, labeled W^+ and W^- .

$$\Rightarrow \langle \sigma_A v \rangle \propto g^4 \sin^2 \theta$$

(b) Equilibrium through coupling to Σ .

A Feynman diagram showing two vertices connected by a cross symbol. The left vertex has two dashed lines labeled S' and Σ' meeting. The right vertex has two dashed lines labeled Σ' and Σ' meeting, with two wavy lines labeled W^+ and W^- emerging from it.

$$\Rightarrow n_{S'}^{-1} \propto \lambda_{\Sigma S'}^2 g^4$$

(c) Equilibrium through coupling to SM Higgs.

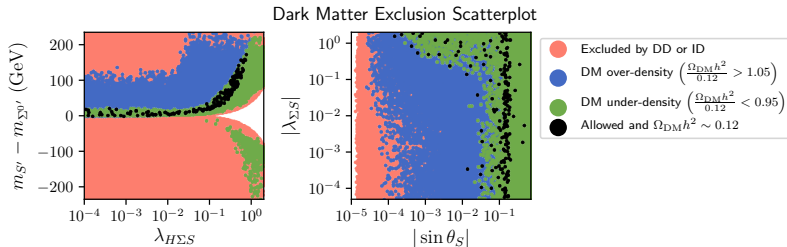
A Feynman diagram showing a vertex where two dashed lines labeled S' meet a dashed line labeled H . From the H line, two solid lines labeled SM emerge.

$$\Rightarrow \langle \sigma_A v \rangle \propto \lambda_{HS'}^2$$

$$\lambda_{HS'} = \lambda_{HS} \cos^2 \theta - \lambda_{H\Sigma S} \cos \theta \sin \theta + \lambda_{H\Sigma} \sin^2 \theta$$

- We use `micrOMEGAS` to evaluate relic density and DD cross sections.
- Apply constraints from XENON1T.
- Apply Fermi-LAT limits arising from loop induced $S'S' \rightarrow \gamma\gamma$ annihilation.
- Also include constraints from oblique parameters.

- Perform a random scan of the parameter space.

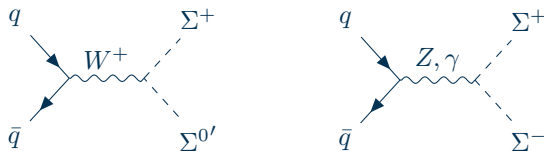


From the scatter-plot, we conclude that either:

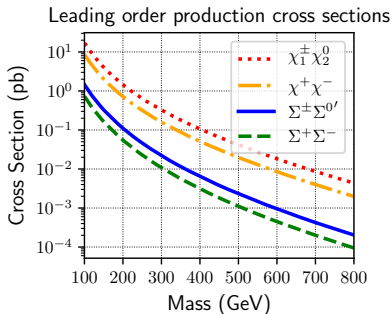
- $|\sin \theta| \sim 0.1$ and $\Delta m > 0$.
- $\lambda_{\Sigma S} > 0.1$ and $0 < \Delta m < 30 \text{ GeV}$.

- No dedicated collider searches exist for this model.
- However, this model is somewhat analogous to SUSY models with a stable neutralino,
 - $\Sigma^\pm \Leftrightarrow \chi_1^\pm$, $\Sigma^{0'} \Leftrightarrow \chi_2^0$, $S' \Leftrightarrow \chi_1^0$
- This model can be constrained using SUSY searches.
- Search strategy depends significantly on “neutralino” mass difference $\Delta m = m_{S'} - m_{\Sigma^{0'}}$.

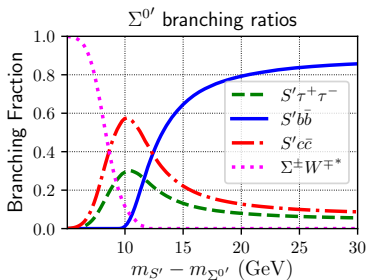
- The scalars are mainly pair produced via charged and neutral current Drell-Yan processes.

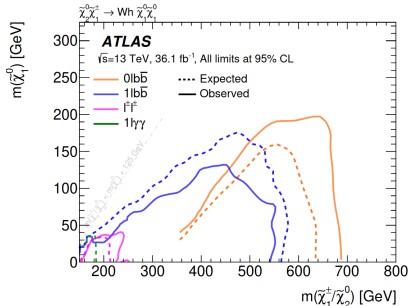
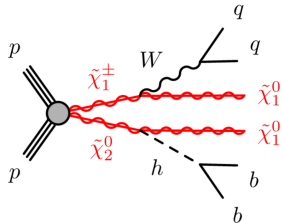


- The production cross section is about $10\times$ smaller than for χ^\pm, χ^0



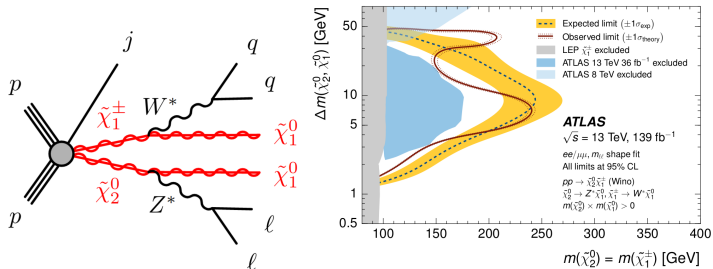
- The Σ^\pm will decay almost exclusively into $S'W^\pm^{(*)}$.
 - Directly analogous to $\chi_1^\pm \rightarrow \chi_1^0 W^\pm^{(*)}$.
- The $\Sigma^{0'}$ will generally decay into $S'h^{(*)}$.
 - For large Δm , there exist analogous neutralino searches using $\chi_2^0 \rightarrow \chi_1^0 h$
 - However, for low Δm , these searches use $\chi_2^0 \rightarrow \chi_1^0 Z^{(*)}$.
- If $m_{\Sigma^{0'}} > m_{\Sigma^\pm}$, the $\Sigma^{0'}$ can also decay into $\Sigma^\pm W^\mp^{(*)}$
 - No direct chargino analogue.





[Image credit: ATLAS Collaboration, arXiv:1812.09432]

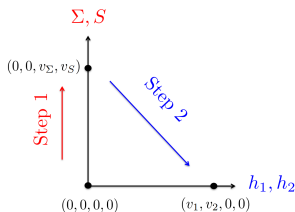
- The large Δm decay processes have directly analogue SUSY searches.
- Constraints well approximated by directly taking the 95% CL on the chargino production cross sections.
- Doing so, we find that current SUSY searches do not constrain our model.



[Image credit: ATLAS Collaboration, arXiv:1911.12606]

- Charginos and neutralinos are notoriously difficult to search for when they only have a small mass difference.
- Existing Searches are not directly applicable due to presence of $Z^* \rightarrow \ell\ell$ decays.
- However, perhaps $h^* \rightarrow \tau\tau \rightarrow \ell\nu\nu\nu\nu$ decays can yield signal events?
- Expect that maybe just about constrain up to $m_{S'} \sim 70$ GeV.

- A Singlet+Triplet extension is capable of explaining the DM density.
- Scalars can be light and have sizeable couplings to SM Higgs \implies potential for novel EWPT.
- Currently relatively unconstrained at colliders.
 - Similar to chargino-neutralino searches, with a factor of $\sim 10\times$ less data.
- Future work
 - How does DM and Collider phenomenology change in Singlet+Triplet+2HDM?
 - This model can generate a baryon asymmetry (arXiv:1508.05404).
 - Can this model simultaneously yield the correct DM density and baryon asymmetry?

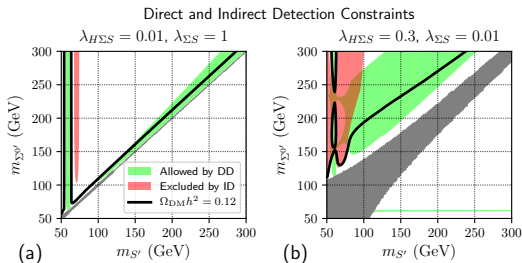
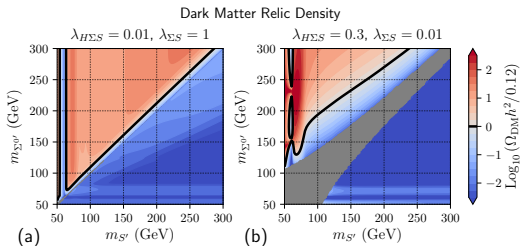


Appendix

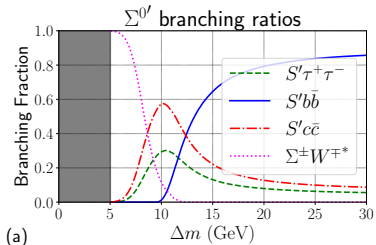
$$\begin{aligned}
V_{\Sigma xSM} = & -\mu_H^2 H^\dagger H - \frac{1}{2}\mu_\Sigma^2 \text{Tr}(\Sigma^2) - \frac{1}{2}\mu_S^2 S^2 \\
& + \lambda_H (H^\dagger H)^2 + \frac{1}{4}\lambda_\Sigma \text{Tr}(\Sigma^2)^2 + \frac{1}{4}\lambda_S S^4 \\
& + \frac{1}{\sqrt{2}}a_{H\Sigma} H^\dagger \Sigma H + a_{HS} H^\dagger H S + \frac{1}{2}a_{\Sigma S} \text{Tr}(\Sigma^2) S + \frac{1}{3}a_S S^3 \\
& + \frac{1}{2}\lambda_{H\Sigma} \text{Tr}(\Sigma^2) H^\dagger H + \frac{1}{2}\lambda_{HS} H^\dagger H S^2 + \frac{1}{4}\lambda_{\Sigma S} \text{Tr}(\Sigma^2) S^2 \\
& + \frac{1}{\sqrt{2}}\lambda_{H\Sigma S} H^\dagger \Sigma H S + bS,
\end{aligned}$$

$$V_{\Sigma xSM} \supset \frac{1}{2} \begin{pmatrix} \Sigma^0 & S \end{pmatrix} \begin{pmatrix} -\mu_\Sigma^2 + \frac{1}{2}v_H^2 \lambda_{H\Sigma} & -\frac{1}{4}v_H^2 \lambda_{H\Sigma S} \\ -\frac{1}{4}v_H^2 \lambda_{H\Sigma S} & -\mu_S^2 + \frac{1}{2}v_H^2 \lambda_{HS} \end{pmatrix} \begin{pmatrix} \Sigma^0 \\ S \end{pmatrix}.$$

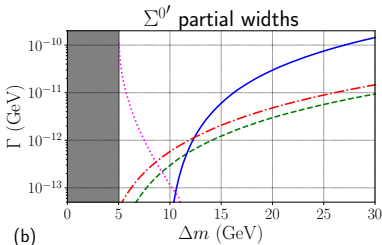
- (a) Equilibrium through weak gauge boson couplings.
- (b) Equilibrium through coupling to Σ .



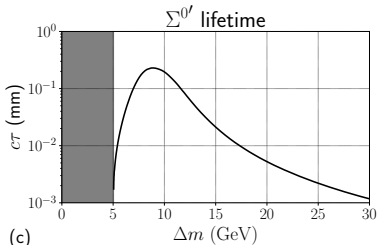
Σ decays with $m_{\Sigma^\pm} = 150$ GeV and $\lambda_{H\Sigma S} = 0.05$



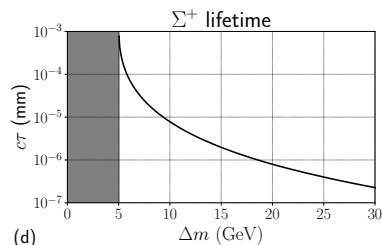
(a)



(b)



(c)



(d)

