



# *Illuminating the Dark Side in Adelaide*

Dipan Sengupta



1. Cosmological constraints on SuperWimps.



Meera Deshpande

+ Hamman, DS, White, Williams, Wong

2. Consequences of neutron decays into dark sector in neutron stars.



Wasif Husain

+ DS, Thomas

+

3. Estimating the relic of KK graviton dark matter accurately.

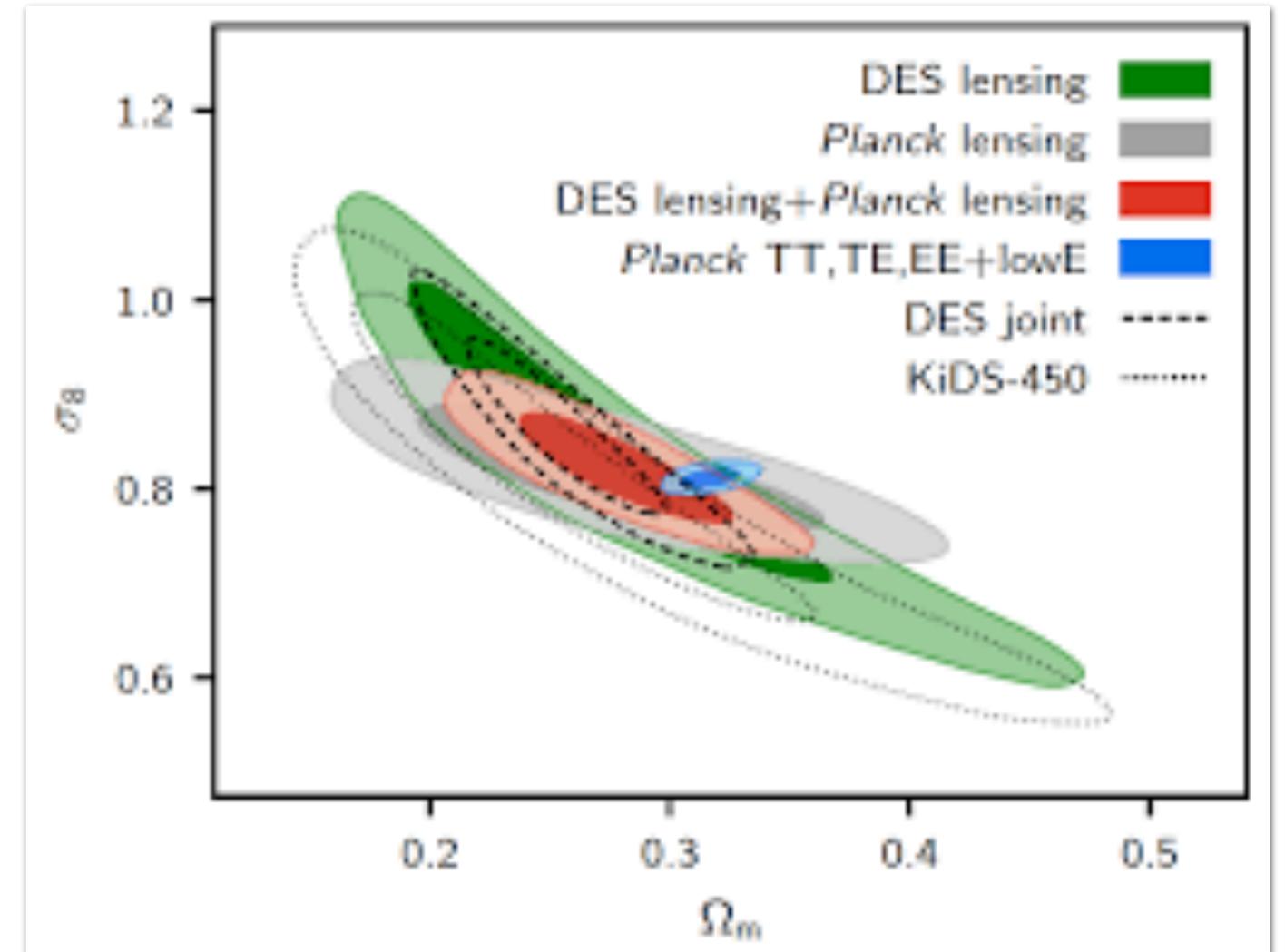
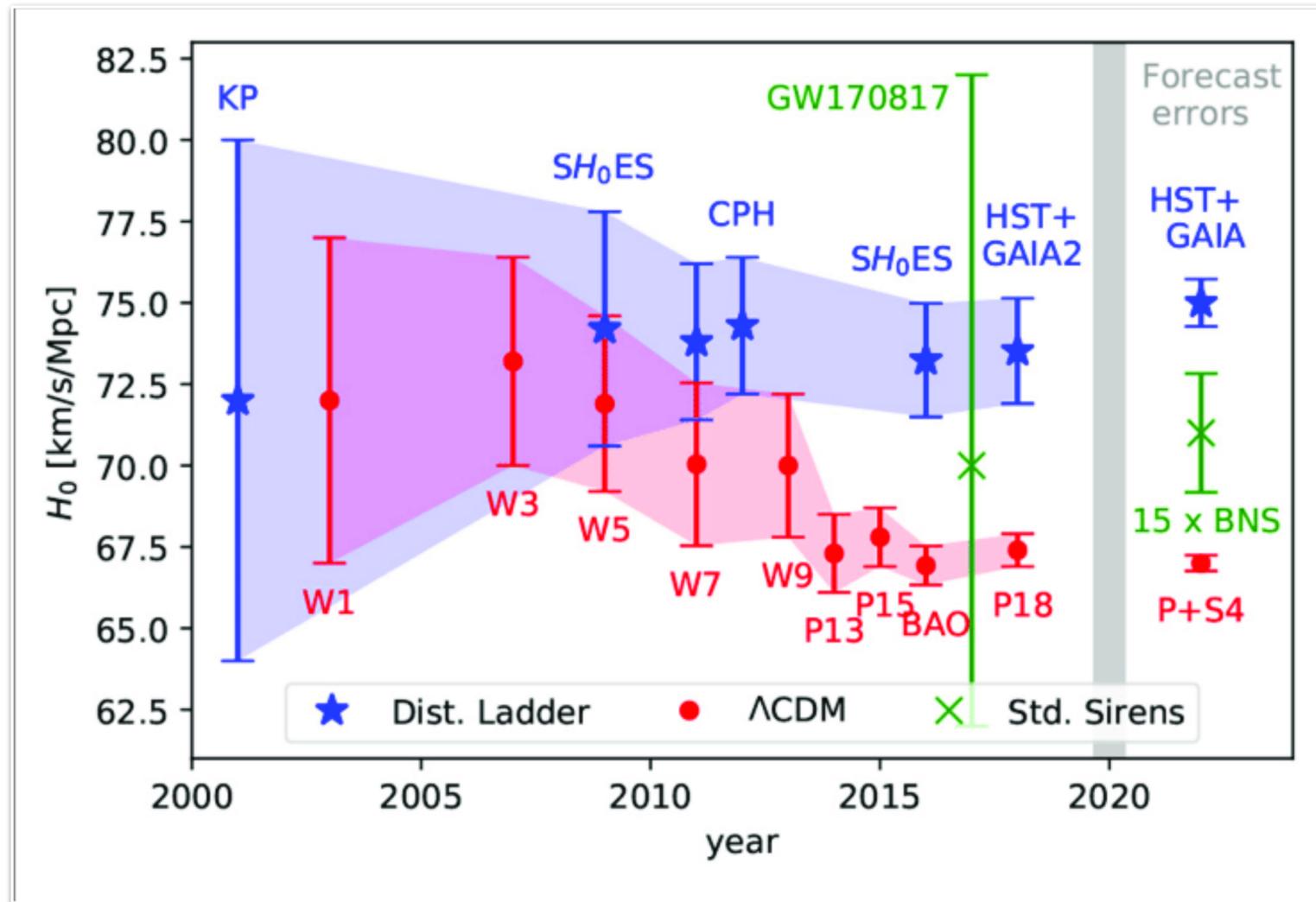


Josh Gill

+ Cacciapaglia(CNRS),  
+ Lee(KIAS), DS, Williams

# Cosmological constraints on SuperWimps.

## Two anomalies in cosmology



$$H_0 = 67.27 \pm 0.60 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

$$H_0 = 74.3 \pm 2.2 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

$$S_8 \equiv \sigma_8 \sqrt{\Omega_m / 0.3}$$

$$S_8 = 0.834 \pm 0.016 \quad \text{Planck}$$

$$S_8 = 0.766^{+0.020}_{-0.014} \quad \text{KiDS-1000}$$

# Cosmological constraints on SuperWimps.

Modifications to  $\Lambda$ CDM model

$S_8$  and the  $H_0$  tensions are correlated

Models of Decaying Dark Matter (DDM) to solve  $S_8$

CDM  $\longrightarrow$  WDM + DR

$$\varepsilon = (1/2) \left[ 1 - m_{\text{wdm}}^2 / m_{\text{dcdm}}^2 \right] \quad \Gamma^{-1} \simeq 55 \text{ Gyrs} \quad \varepsilon \simeq 0.7 \% \quad \text{Poulin-Abellan-Lavalle-Murgia : 2020}$$

Suppression of Linear Matter power spectrum at intermediate and small scales with a cut-off scale determined by the free streaming length

Difference between  $\Lambda$ CDM and  $\Lambda$ DDM : Very small at low redshifts , and therefore Planck cannot distinguish them

# Cosmological constraints on SuperWimps.

What kind of models can we construct? Look no further than SUSY

Yanagagida et al. 2020

Consider a gravitino CDM populated thermally in the early universe through scatterings

$$\Omega_{3/2} h^2 = 0.217 \left( \frac{T_{\text{RH}}}{10^7 \text{GeV}} \right) \left( \frac{100 \text{GeV}}{m_{3/2}} \right) \left( \frac{m_{\tilde{g}}(\mu)}{10 \text{TeV}} \right)^2$$

$$m_{3/2} = \frac{|F|}{\sqrt{3} M_P}$$

$$\tilde{G}_\mu \rightarrow \tilde{N}_1 + N_1$$

$$\Gamma(\tilde{G}_\mu \rightarrow \tilde{N}_1 + N_1) = \frac{m_{3/2}^3}{192\pi M_P^2} \times \left[ 1 - \left( \frac{m_1}{m_{3/2}} \right) \right]^2 \left[ 1 - \left( \frac{m_1}{m_{3/2}} \right)^2 \right]^3$$

Solves the Sigma\_8 tension

# Cosmological constraints on SuperWimps.

What if the reheating temperature is low ? Thermal processes are suppressed

Gravitino abundance is populated non thermally through decays

$$\Gamma(\chi_1^0 \rightarrow \tilde{G}\gamma) \equiv \frac{\cos^2 \theta_W m_{\chi_1^0}^5}{48 M_P^2 m_{\tilde{G}}^2} \left[ 1 - \frac{m_{\tilde{G}}^2}{m_{\chi_1^0}^2} \right]^3 \left( 1 + 3 \frac{m_{\tilde{G}}^2}{m_{\chi_1^0}^2} \right)$$

$$\tau \equiv 2.3 \times 10^7 \left( \frac{100 \text{ GeV}}{\Delta m} \right)^3 \text{ s}$$

Energy released in Photons

$$E_\gamma = \frac{m_{\chi_1^0}^2 - m_{\tilde{G}}^2}{2m_{\chi_1^0}}$$

Fractional energy

$$E_{\text{SM}} = E_\gamma / m_{\chi_1^0}$$

**Energy deposited in the thermal plasma causes spectral distortions**

# Cosmological constraints on SuperWimps.

## Energy Injection Constraints

## Spectral Distortions

### Distortions of the Blackbody spectrum of the primordial photon bath

Energy injection and deposition into the Intergalactic Medium (IGM)

$$\left. \frac{dE}{dt dV} \right|_{\text{dep,c}} = \left. \frac{dE}{dt dV} \right|_{\text{inj}} f_c = \left. \frac{dE}{dt dV} \right|_{\text{inj}} f_{\text{eff}} \chi_c \equiv \dot{Q} \chi_c$$

injection efficiency function  $f_{\text{eff}}(z)$   
deposition fraction  $\chi_c(z)$

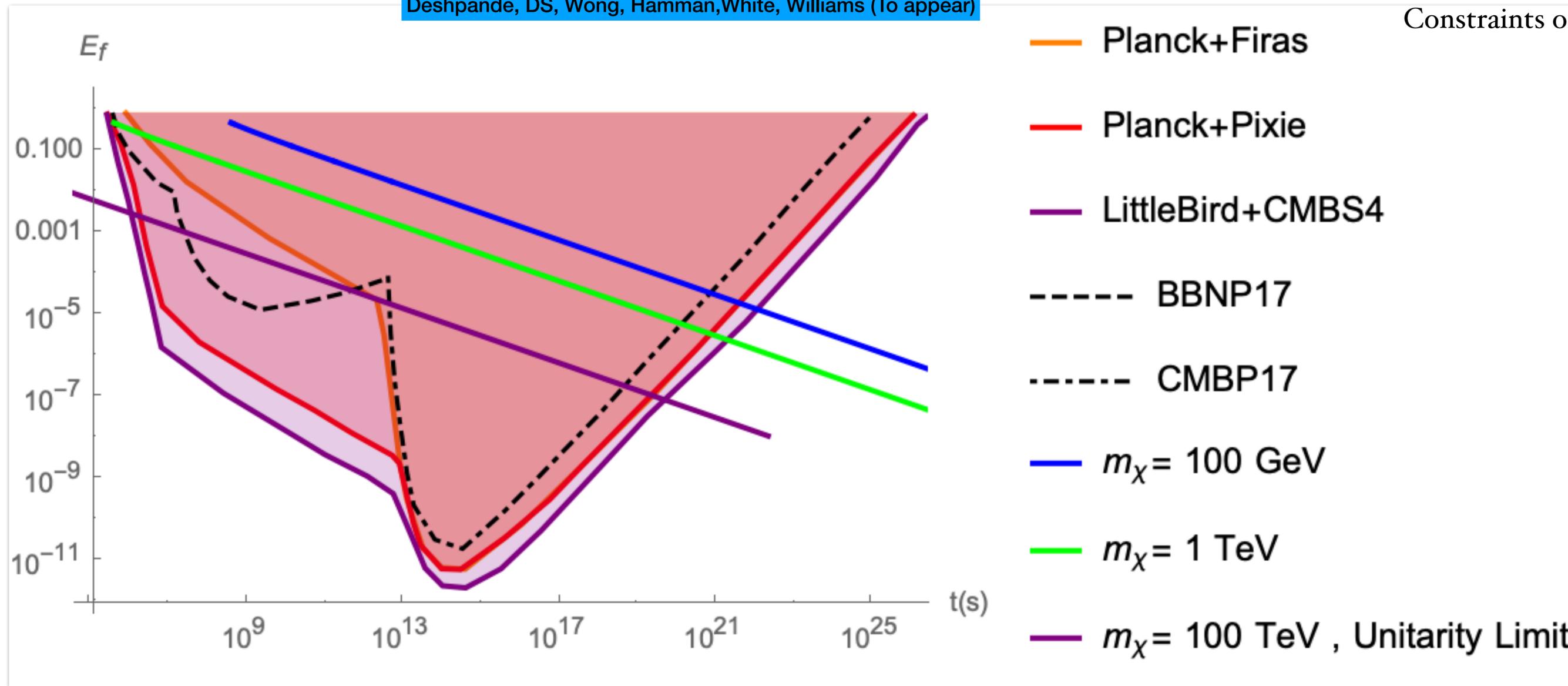
Thermalisation via Compton, Double Compton and Brehmstrahlung scatterings  $\longrightarrow$  Photon Phase Space Distribution

Distortions manifested in terms of temperature shifts  $\mathbf{g}$ , chemical potential distortions  $\mathbf{mu}$ , and Compton distortions  $\mathbf{y}$

# Cosmological constraints on SuperWimps.

Deshpande, DS, Wong, Hamman, White, Williams (To appear)

Constraints on Gravitino SuperWIMP



**Similar considerations for axino SuperWimps : Additional freedom in decay width due to axion decay constant**

In consideration : Complementarity between collider, Warm DM bounds  
 Future : Axino/Gravitino decays for solving Hubble/ $S_8$  tensions consistent with constraints  
 Release Code to do understand general multistep process in Class/Exoclass  
 Other DDM scenarios

1. Cosmological constraints on SuperWimps.



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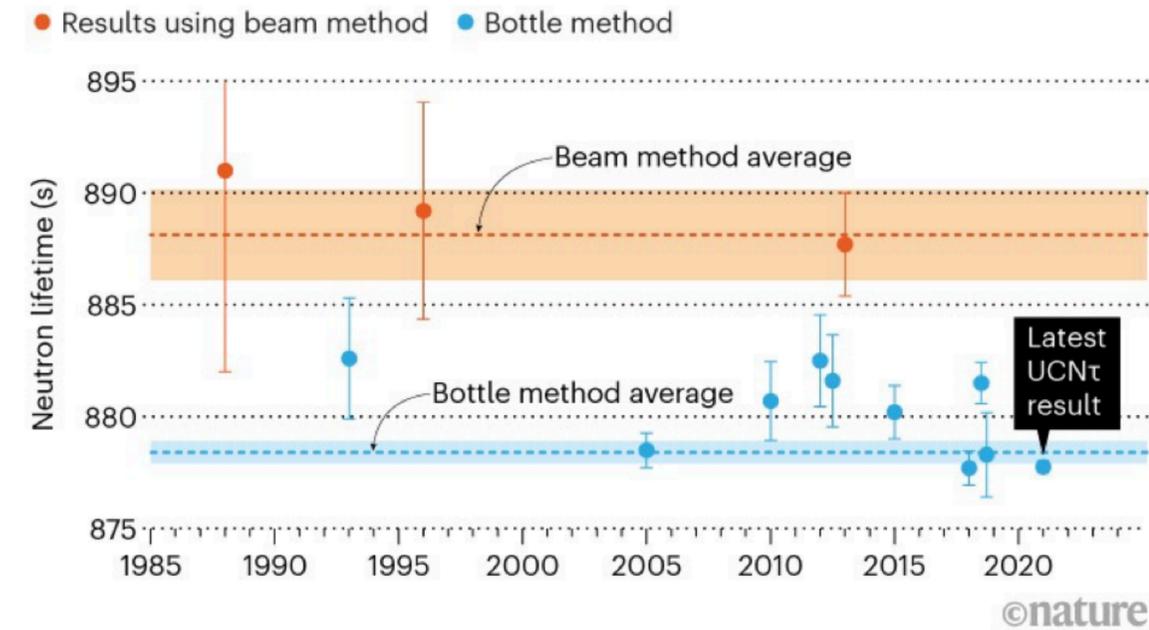
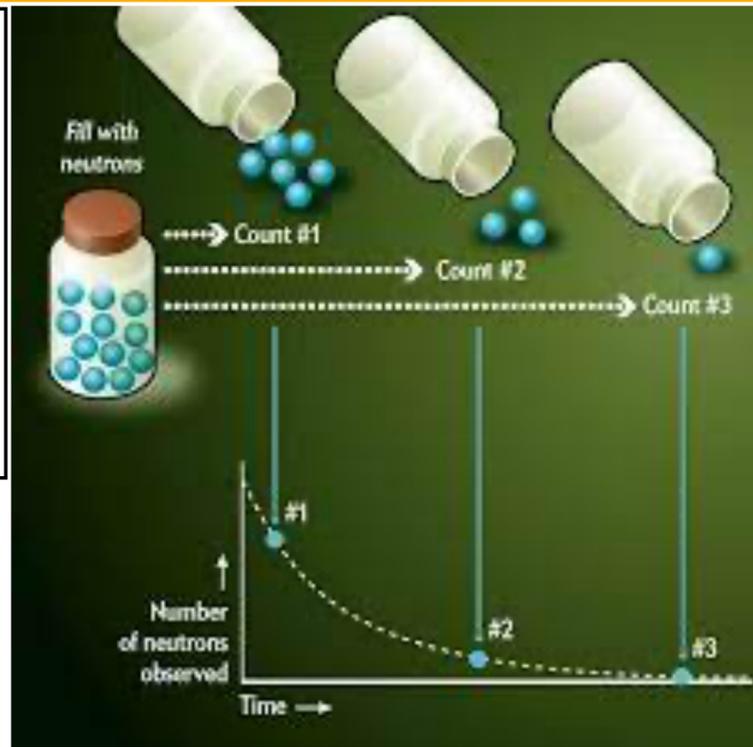
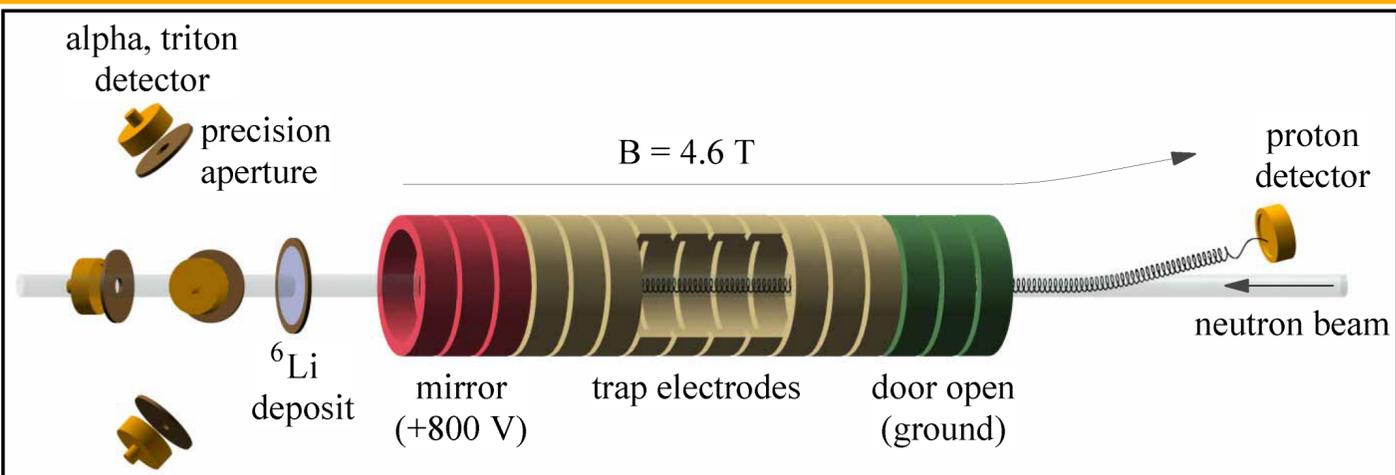
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# Bottle vs Beam experiments



$$\mathcal{M} = \frac{1}{\sqrt{2}} G_F V_{ud} g_V [\bar{p} \gamma_\mu n - \lambda \bar{p} \gamma_5 \gamma_\mu n] [\bar{e} \gamma^\mu (1 - \gamma_5) \nu]$$

$$\tau_n = \frac{4908.7(1.9) \text{ s}}{|V_{ud}|^2 (1 + 3\lambda^2)}$$

$\tau_n$  between 875.3 s and 891.2 s within  $3\sigma$

$$\tau_n^{\text{beam}} = 888.0 \pm 2.0 \text{ s}$$

$$\tau_n^{\text{bottle}} = 879.6 \pm 0.6 \text{ s}$$

$$\Delta\Gamma_n^{\text{exp}} = \Gamma_n^{\text{bottle}} - \Gamma_n^{\text{beam}} \simeq 7.1 \times 10^{-30} \text{ GeV}$$

# New Physics Interpretations

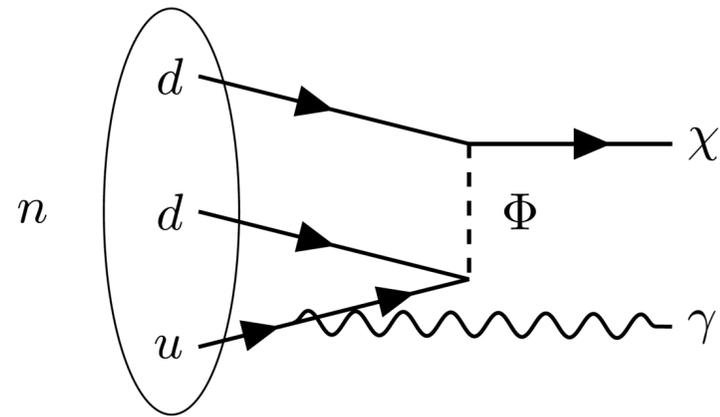
## New Physics scenarios :

### Set I

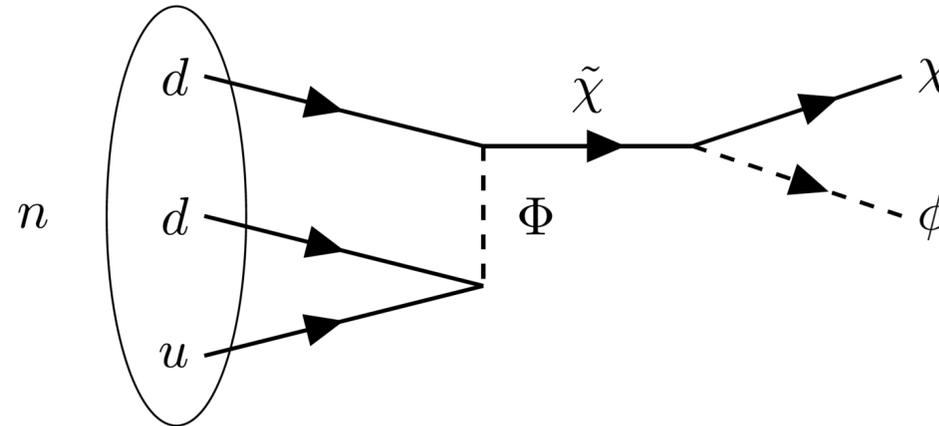
$$n \rightarrow \chi \gamma$$

$$n \rightarrow \chi \phi$$

$$n \rightarrow \chi e^+ e^-$$



Fornal-Grienstein, Nelson et al ...



$$\Phi = (3, 1)_{-1/3}$$

$$937.900 \text{ MeV} < m_\chi + m_\phi < 939.565 \text{ MeV}$$

### Set II

DM quantum numbers			DM interactions		
$B$	$L$	spin	dimension	with quarks	with hadrons
1	0	1/2	6	$\chi u d d$	$\chi n$
1/3	0	1/2	9	$\chi \chi \chi u d d$	$\chi \chi \chi n$
2/3	0	0	9	$\phi^3 (u d d)^2$	$\phi^3 n^2$
2	0	0	7	$\phi (u d d)^2$	$\phi n n$
0	1	1/2	4, 6	$\chi L H, \chi l f \bar{f}$	$\chi l \pi, \chi l p \bar{n}$
0	2	0	6, 8	$\phi (L H)^2, \phi l l X q \bar{q}$	$\phi \nu \nu, \phi l l \pi \pi$
1	1	0	7	$\phi L Q Q Q, \phi l u u d$	$\phi n \nu, \phi p l$
1	-1	0	8	$\phi \bar{l} X q q q$	$\phi n \bar{\nu}, \phi \Delta^- \bar{l}, \phi n \pi^- \bar{l}$
1	2	1/2	9	$\chi l \nu q q q$	$\chi n \nu \nu, \chi p l \nu$

Strumia's classification

$$n \rightarrow \chi \chi \chi$$

# Neutron Star Considerations with the decay

The new decay softens the neutron star EOS at high densities  $\longrightarrow$  Makes it impossible to support NS above 2 solar masses

## Two Solutions

1. Large repulsive self interactions between DM, stiffens the EOS by raising DM chemical potential, reduces DM to baryon fraction in equilibrium
2. Repulsive DM-Baryon interactions : energetically disfavours DM production in a pure baryonic medium

$$U = \pm \frac{g_\chi g_n}{4\pi} \frac{e^{-m_\phi r}}{r}$$

*TOV equation for hydrostatic equilibrium with DM and Neutrons*

Grinstein et al. 2018

$\phi$  In principle can cause problems by adding to  $N_{\text{eff}}$  : Ideally should decay before start of BBN to avoid all constraints

# New Physics Interpretations

## What if the boson decayed ?

$$n \rightarrow \chi \phi \quad 937.900 \text{ MeV} < m_\chi + m_\phi < 939.565 \text{ MeV}$$

**Both Stable if**  $|m_\chi - m_\phi| < 938.783 \text{ MeV}$

## How do we estimate the mass and the lifetime of the Boson ?

Urca and inverse Urca processes cool neutron stars down

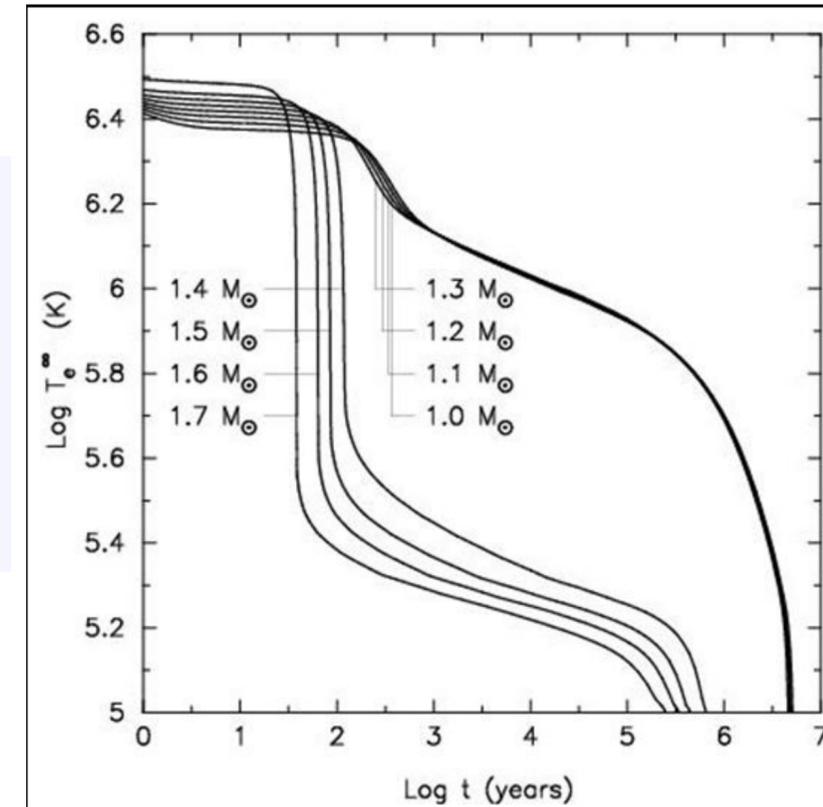
neutron stars have a luminosity  $10^{31.5}$  erg/s at 1 M years

Cooling Processes	
➤ Direct Urca:	$n \rightarrow p + e + \bar{\nu}_e$
➤ Modified Urca:	$n + n \rightarrow n + p + e + \bar{\nu}_e$
➤ Photons:	$\rightarrow \gamma$
➤ Bremsstrahlung:	$n + n \rightarrow n + n + \nu + \bar{\nu}$

Additional cooling process due to SM particles from boson decays should not cool it below  $10^{31.5}$  erg/s

$m_\phi$     **300 MeV**     **$47 \times 10^{13}$  years**

**700 MeV**     **$70 \times 10^{13}$  years**



# New Physics Interpretations

## Scalars

$$\mathcal{L}_{int} = \frac{C_s}{f_{eff}} \phi F_{\mu\nu} F^{\mu\nu} + \frac{m_l}{f_{eff}} \phi \bar{l} l + \dots$$

$$\mathcal{L} \in L_{kin} + \lambda_{eff} n \chi \phi$$

## Spin-1

Photon is (almost) ruled out experimentally

$$\mathcal{L} = \frac{\epsilon}{2 \cos \theta_W} \tilde{F}'_{\mu\nu} B^{\mu\nu}$$

$$\mathcal{L} \in e\epsilon(n\sigma^{\mu\nu}\chi F'_{\mu\nu})$$

## Pseudo-Scalars

$$\mathcal{L}_{int} = \frac{C_{s\gamma}}{f_{eff}} \phi F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{f_{eff}} (\partial_\mu \phi) \bar{l} \gamma^\mu \gamma^5 l + \dots$$

$$\mathcal{L} \in L_{kin} + \lambda_{eff} n \chi \gamma^5 \phi$$

## Spin-2

$$\mathcal{L} = \frac{1}{\Lambda} h_{\mu\nu} T_{SM}^{\mu\nu}$$

In the works : A full analysis of complementary constraints on this scenario including stellar cooling bounds, low energy and collider experiments, BBN and CMB bounds

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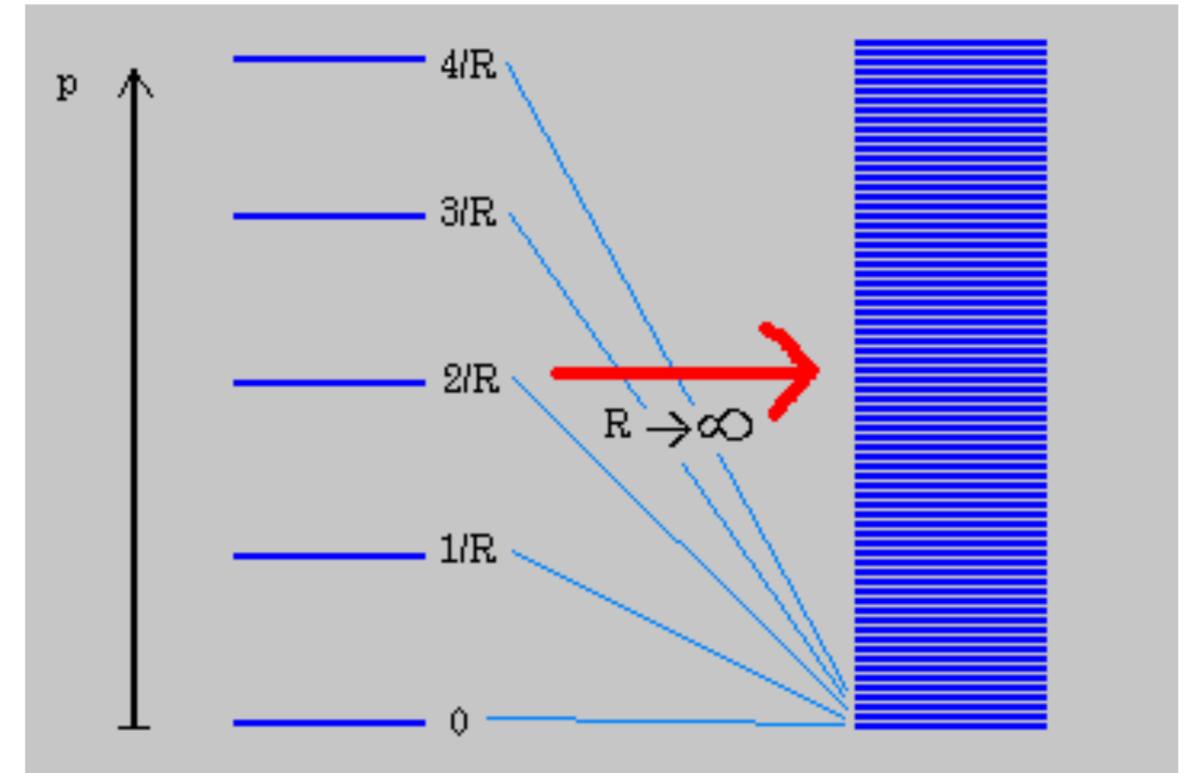
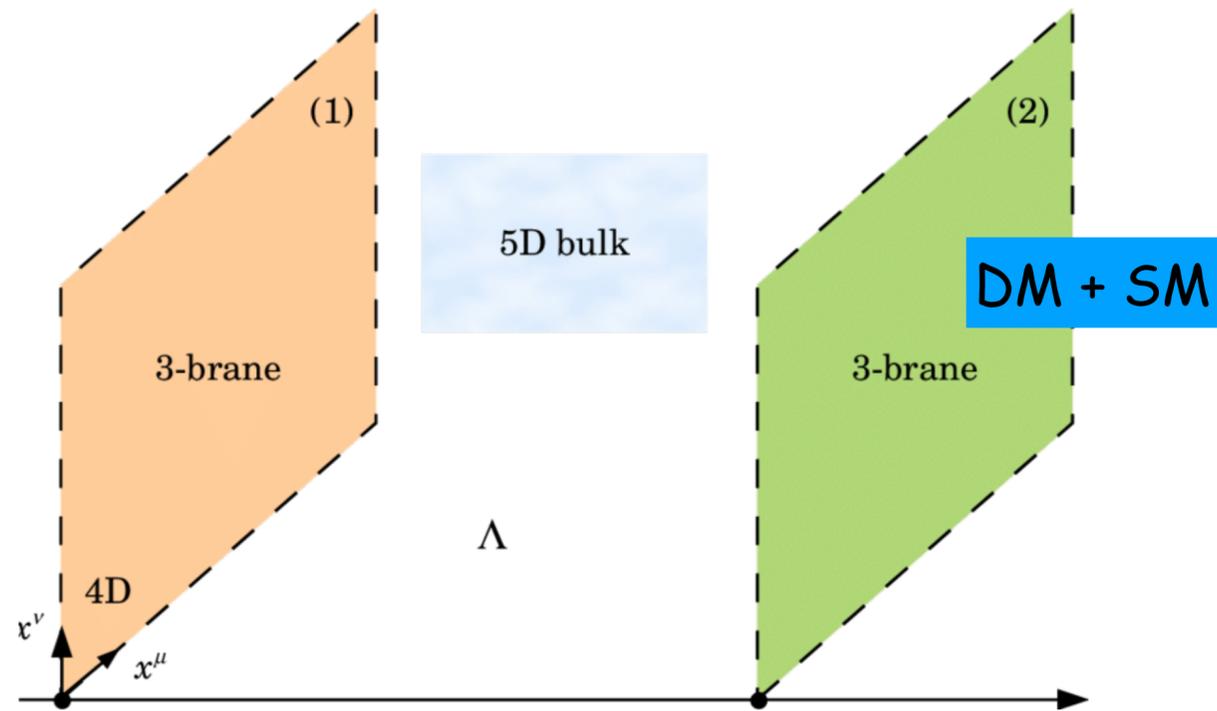


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# Estimating Accurate Relic Densities for DM models with KK gravitons

Massive Spin-2 KK gravitons arise as a result of compactifying extra dimensions



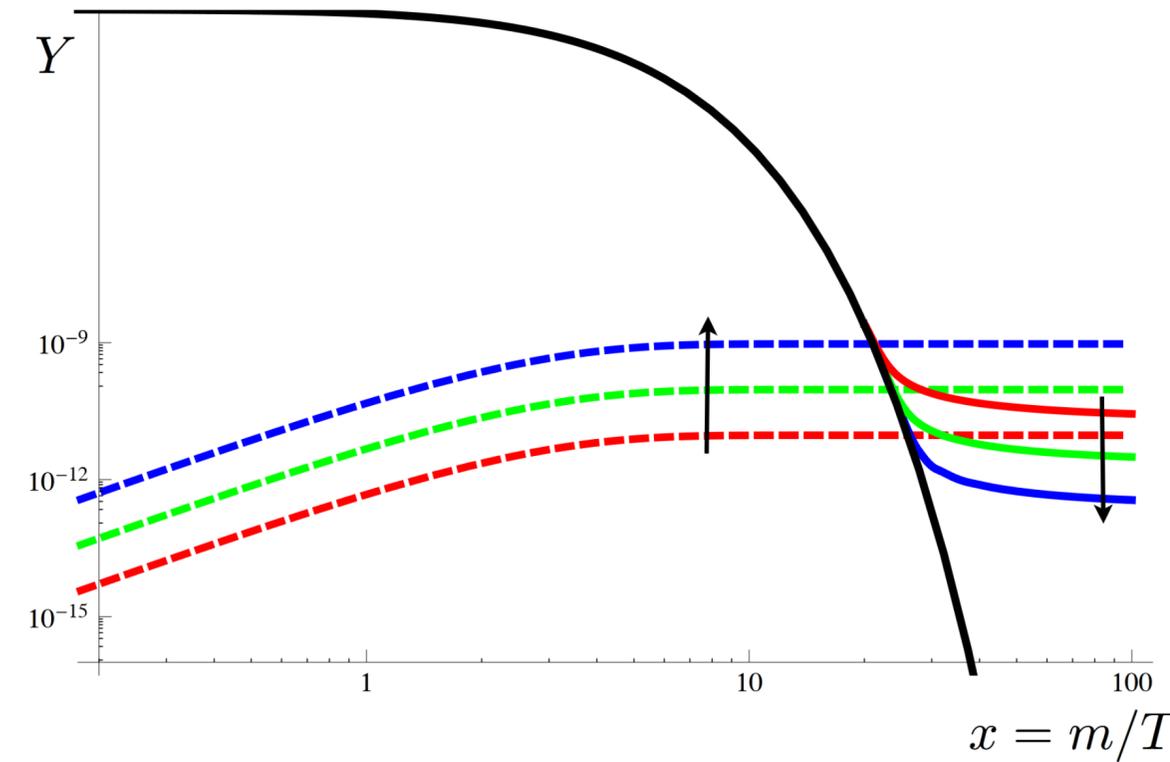
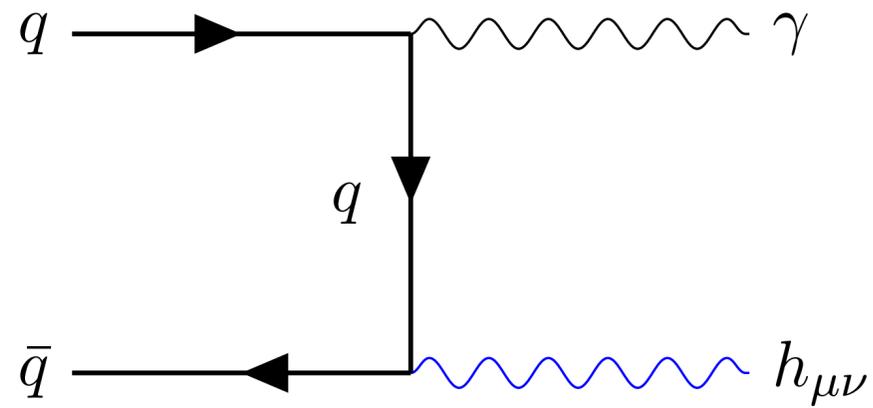
$$\mathcal{L} = \frac{1}{\Lambda} h_{\mu\nu} T_{SM}^{\mu\nu}$$

**Naive Expectation of all EFT scales in the theory**

# Estimating Accurate Relic Densities for DM models with KK gravitons

A light KK graviton with a lifetime greater than the age of the Universe.

UV freeze-in through higher dimensional operators

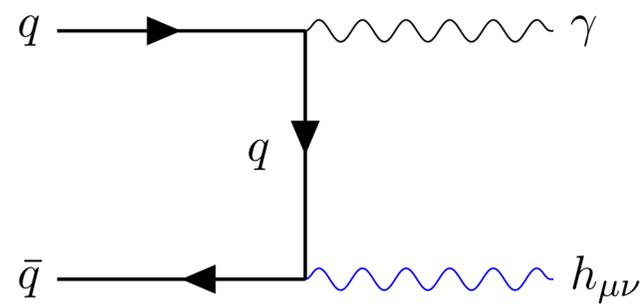


$$\dot{n}_\chi + 3Hn_\chi = \sum_{A,B} (\xi_B - \xi_A) \mathcal{N}(A \rightarrow B)$$

The matrix element contains information about the velocity averaged cross section

Unitarity limits for effective theories determine the validity of the theory

# Estimating Accurate Relic Densities for DM models with KK gravitons



+t + u + contact diagrams

$$\lambda_G = \pm 2, \quad \varepsilon_{\pm 2}^{\mu\nu} = \varepsilon_{\pm 1}^{\mu} \varepsilon_{\pm 1}^{\nu},$$

$$\lambda_G = \pm 1, \quad \varepsilon_{\pm 1}^{\mu\nu} = \frac{1}{\sqrt{2}} \left[ \varepsilon_{\pm 1}^{\mu} \varepsilon_0^{\nu} + \varepsilon_0^{\mu} \varepsilon_{\pm 1}^{\nu} \right],$$

$$1: \lambda_G = 0, \quad \varepsilon_0^{\mu\nu} = \frac{1}{\sqrt{6}} \left[ \varepsilon_{+1}^{\mu} \varepsilon_{-1}^{\nu} + \varepsilon_{-1}^{\mu} \varepsilon_{+1}^{\nu} + 2\varepsilon_0^{\mu} \varepsilon_0^{\nu} \right]$$

$$\varepsilon_0^{\mu} (k_2) = \frac{E_{k_2}}{m_G} \left( \sqrt{1 - \frac{m_G^2}{E_{k_2}^2}}, \hat{k} \right)$$

Matrix Element naively grows like  $1/M_{\text{KK}}^2$

**Only one EFT scale, should not have any low energy divergences**

Incorrect estimations in the literature , Lee et al, Sanz et al,  
Sloth et al, Bernal et al, Mambrini et al ....

**Solution : Sum the KK tower , All low energy divergences should cancel out**

**Status : Messy matrix elements, 40 Helicity combinations, manipulations to get them into a tractable form.  
Sum the KK tower, Calculate the cross section , integrate the Boltzmann Equation**

# Summary

The Adelaide Theory Group has a rich DM/Cosmology theory programme working on a variety of topics

A lot of scope and directions to collaborate

