

Imprints of the Early Universe on Dark Matter Substructure

Matthew Dolan
University of Melbourne
Centre of Excellence for Dark Matter Particle Physics

Based on 1905.06952 and 1911.07853 with Nikita Blinov, Jonathan Kozaczuk and Patrick Draper.



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Axion-Like Particles

- Axions/ALPs: (spin 0) Pseudo-Nambu-Goldstone bosons, or zero-modes of higher dimensional gauge fields.
- Naturally light, (very) weakly coupled. Well-known dark matter candidate.
- Will not discuss QCD axion (solution of Strong CP problem)

$$\mathcal{L} \supset -\frac{1}{2}m_a^2 a^2 + \frac{1}{4}g_{a\gamma\gamma}F\tilde{F}$$
$$g_{a\gamma\gamma} \sim \frac{\alpha}{2\pi f_a}$$

- QCD axion has $m_a \sim m_\pi f_\pi / f_a$
- ALP has m_a and f_a independent.
- Also other possible couplings.

ALP Dark Matter: Misalignment

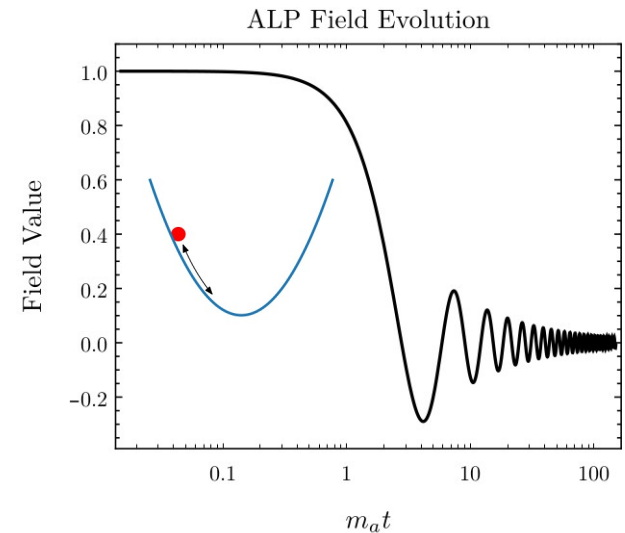
- The axion/ALP is not a thermal relic (i.e. not a WIMP)
- Field displaced from origin in early universe
- Starts to oscillate around origin when

$$m_a \sim H(T_{\text{osc}})$$

- The ALP field evolves as

$$a \sim a_i \left(\frac{R_0}{R(t)} \right)^{3/2} \cos m_a t$$

- $R(t)$ is the scale factor of the FRW cosmology.



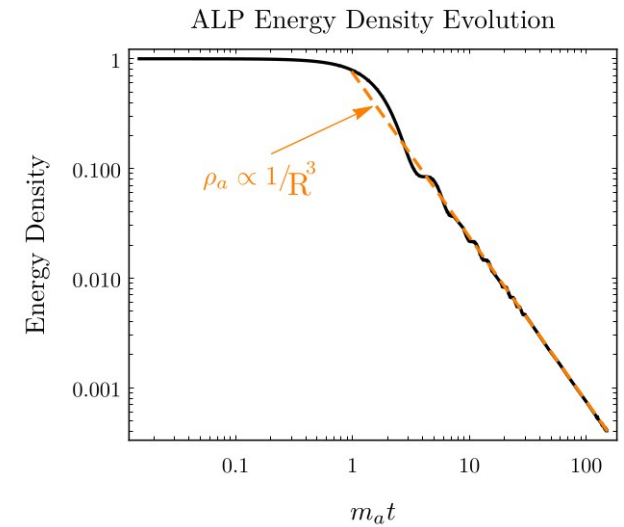
ALP Dark Matter: Misalignment

- Average over oscillation periods, find $\rho_a \sim 1/R^3$
- ALPs behave as CDM (mostly).
- Assuming standard cosmology with radiation domination in early universe:

$$\Omega_a = \frac{1}{2} \frac{m_a^2 f_a^2 \theta_0^2}{\rho_c} \left(\frac{R_{\text{osc}}}{R_0} \right)^3$$

ALP abundance depends on initial misalignment

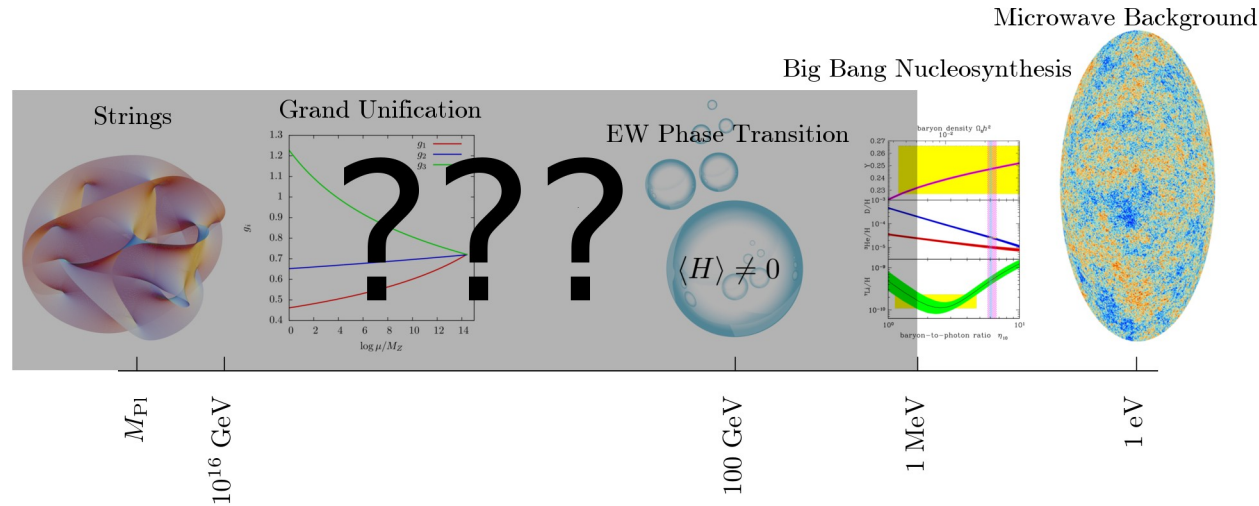
ALP abundance depends on expansion since oscillations began.



$$\Omega_a h^2 \simeq 0.12 \left(\frac{f_a \theta_0}{1.9 \times 10^{13} \text{ GeV}} \right)^2 \left(\frac{m_a}{1 \text{ } \mu\text{eV}} \right)^{1/2}$$

ALP Dark Matter: Misalignment

- Don't know the evolution of the universe before BBN.

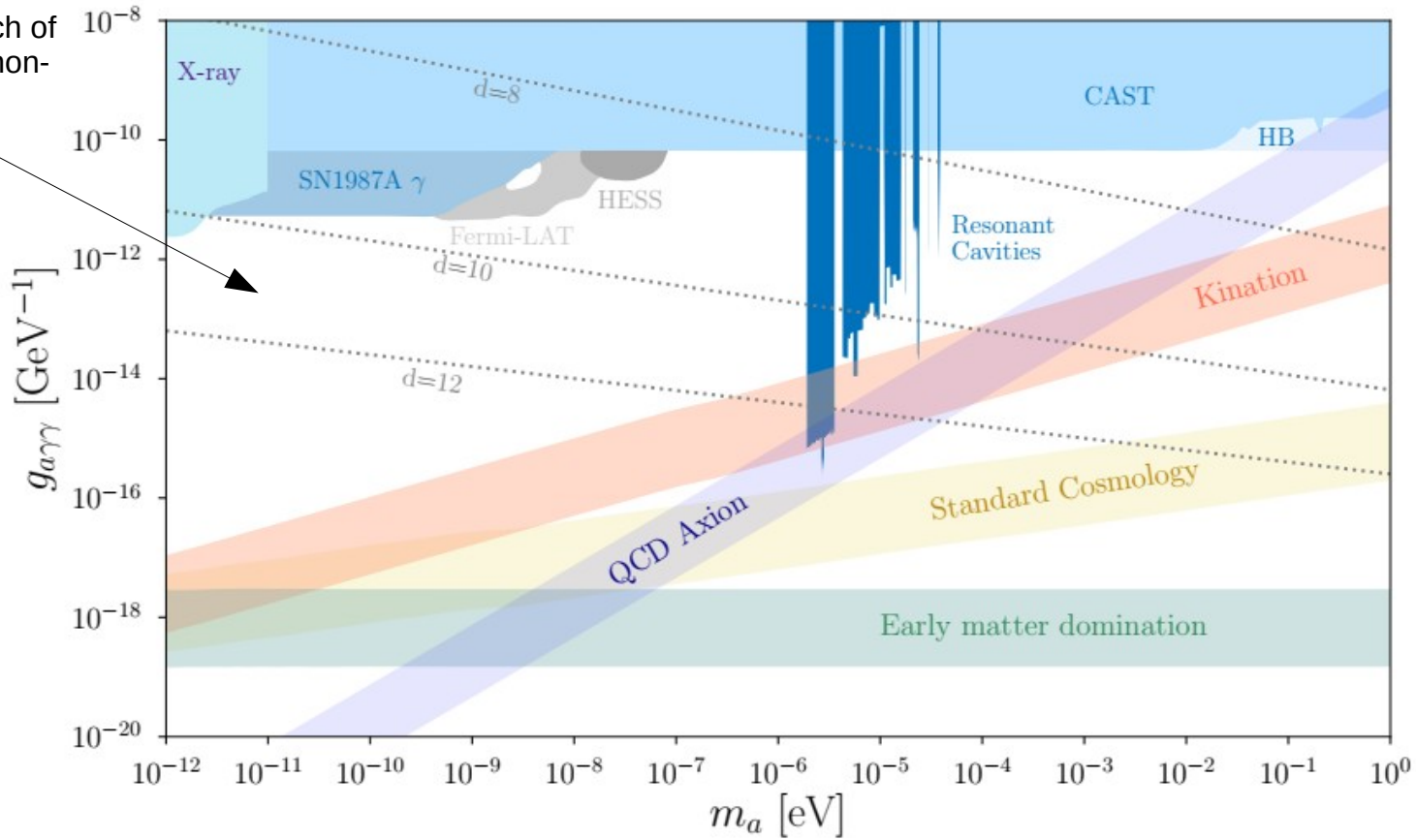


- ALP UV completions often involve an epoch of Early Matter Domination (EMD).
- Transition from EMD to radiation domination (reheating – RH) happens at T_{RH}
- BBN constraint: $T_{RH} > 5$ MeV

$$\Omega_a h^2 \simeq 0.12 \left(\frac{f_a \theta_0}{9 \times 10^{14} \text{ GeV}} \right)^2 \left(\frac{T_{RH}}{10 \text{ MeV}} \right)$$

Current Constraints

Can also fill this much of this region out with non-cosmological modifications

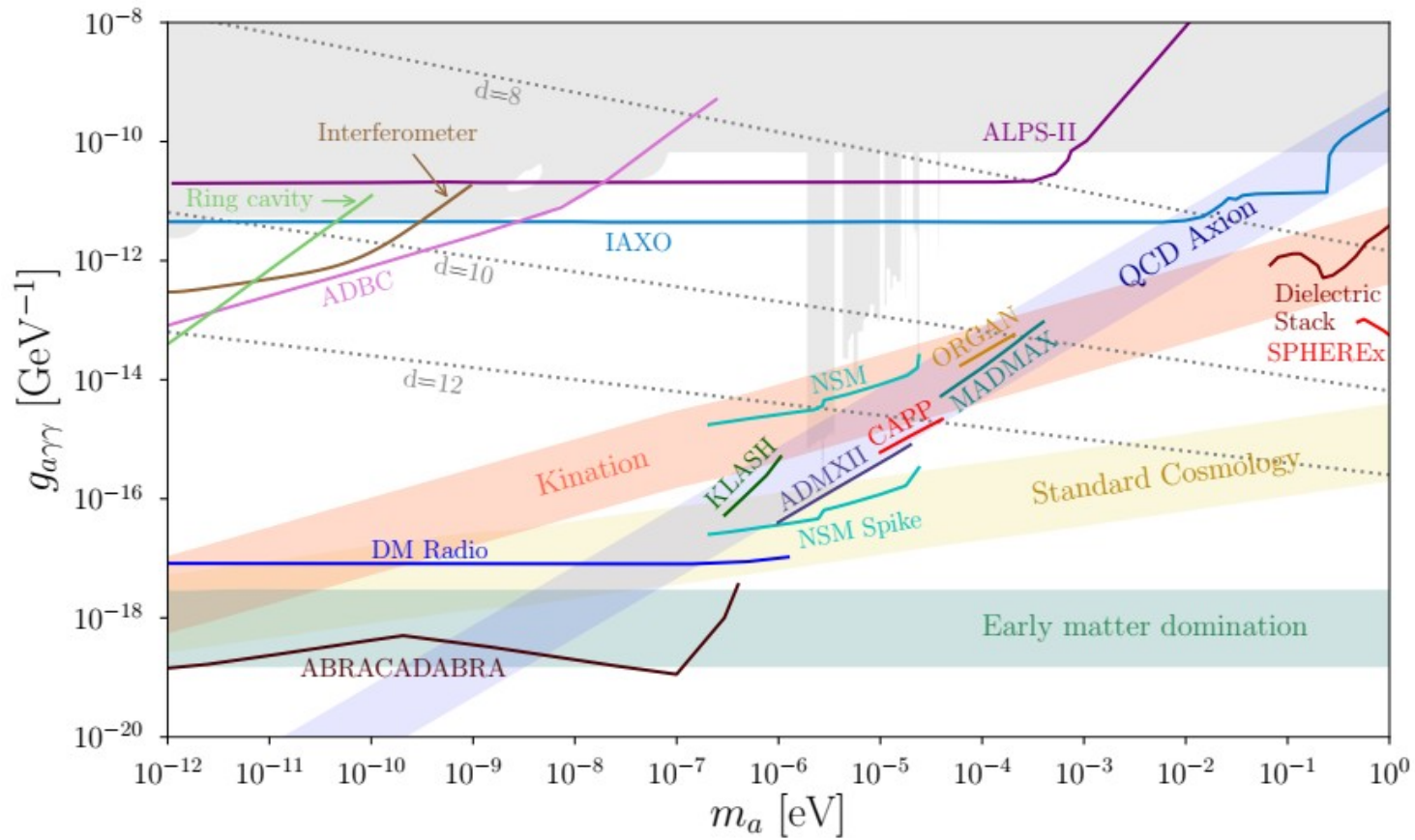


Non-DM out here: collider, astrophysical, flavour constraints.

MJD, Kahlhoefer, McCabe, Schmidt-Hoberg 2014;
MJD, Ferber, Hearty, Kahlhoefer, Schmidt-Hoberg 2017

MJD, Hiskens, Volkas, to appear

Current Constraints

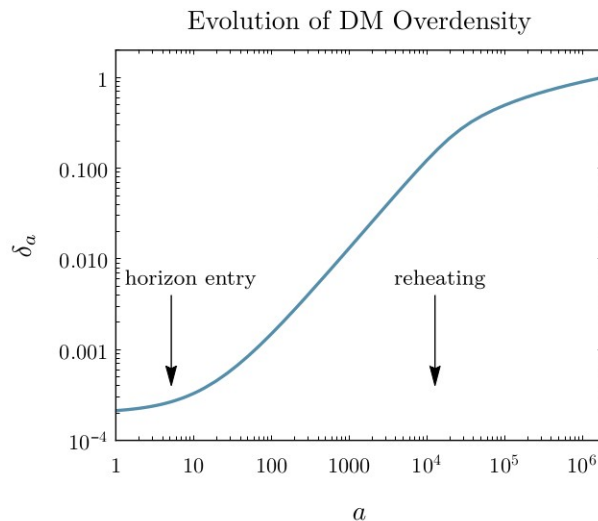


Impact on Small Scale Structure: Clumps

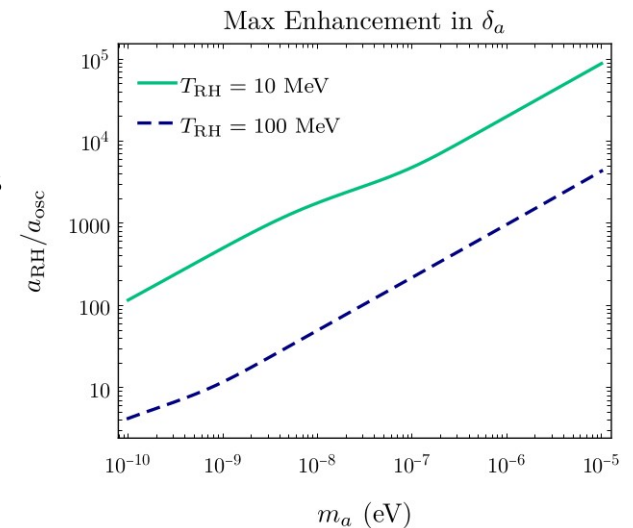
- Changing cosmology changes growth of density perturbations in early universe.
- Consider a CDM density perturbation $\delta = (\rho(x) - \bar{\rho}) / \bar{\rho}$
- Evolves as $\ddot{\delta} + \mathcal{H}\dot{\delta} \approx -k^2\Psi$ where k = comoving inverse size of perturbation

$$\delta \propto \begin{cases} a & \text{MD} \\ \ln a & \text{RD} \end{cases}$$

- EMD enhances growth by $\sim (a_{\text{RH}}/a_{\text{hor}}) \sim (k/k_{\text{RH}})^2$

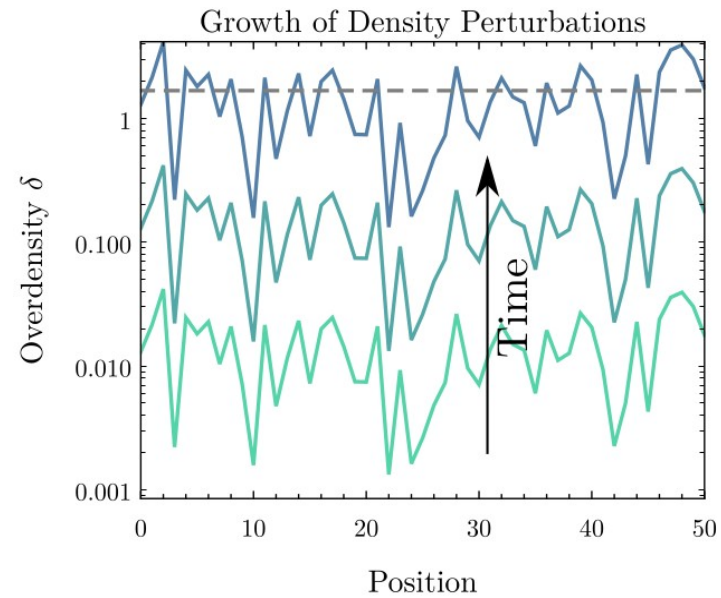


$$\frac{a_{\text{RM}}}{a_{\text{OSC}}} \sim 10^6 \left(\frac{m_a}{10^{-5} \text{ eV}} \right)^{2/3} \left(\frac{5 \text{ MeV}}{T_{\text{RH}}} \right)^{4/3}$$



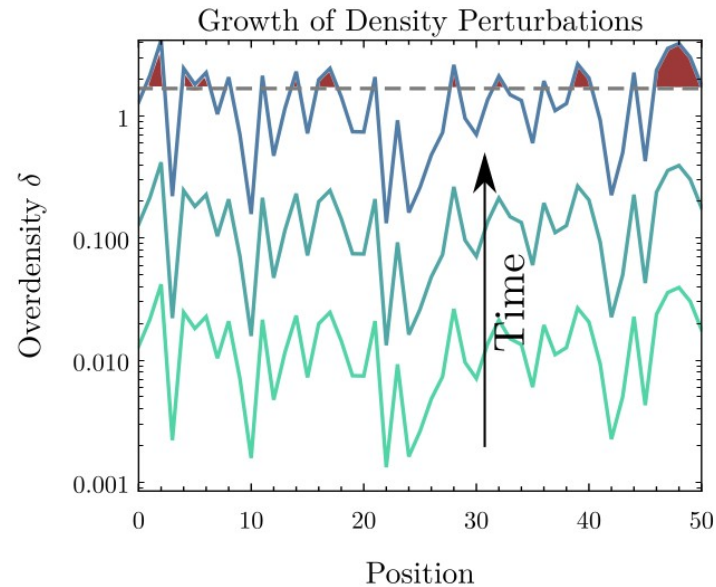
Typical Clumps

- The early universe is a realisation of the primordial density distribution



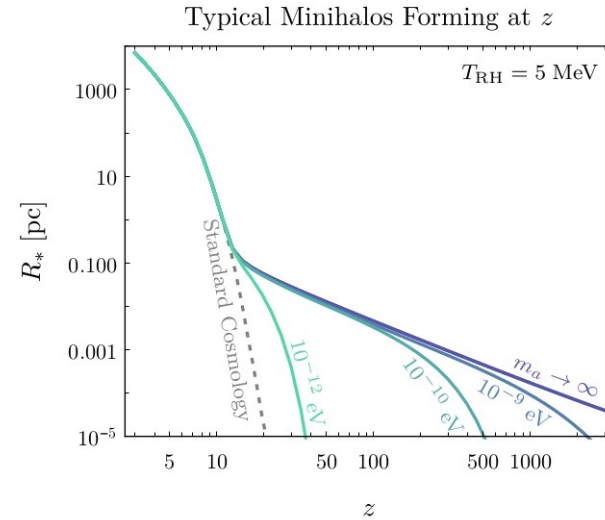
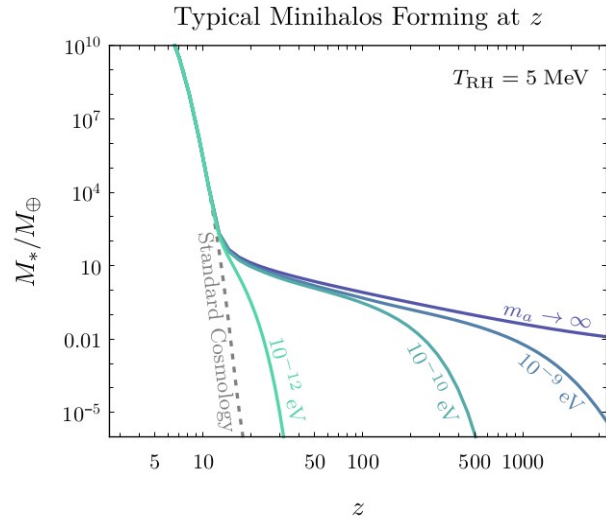
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- If a clump grows dense enough, decouples from Hubble flow and collapses at redshift z_c and forms a virialised minihalo.
- Study statistics of collapsed density perturbations using Press-Schechter theory.

Typical Clumps



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$$\rho(z_c) \approx 230 \text{ GeV/cm}^3 \left(\frac{1+z_c}{100} \right)^3$$

$$M_*(z_c) \approx 250 M_\oplus \left(\frac{5 \text{ MeV}}{T_{\text{RH}}} \right)^3 \left(\frac{2.6}{1+z_c} \right)^{6/(n_s+3)}$$

$$R_*(z_c) \approx 4 \times 10^{-3} \text{ pc} \left(\frac{5 \text{ MeV}}{T_{\text{RH}}} \right) \left(\frac{100}{1+z_c} \right)^{\frac{5+n_s}{3+n_s}}$$

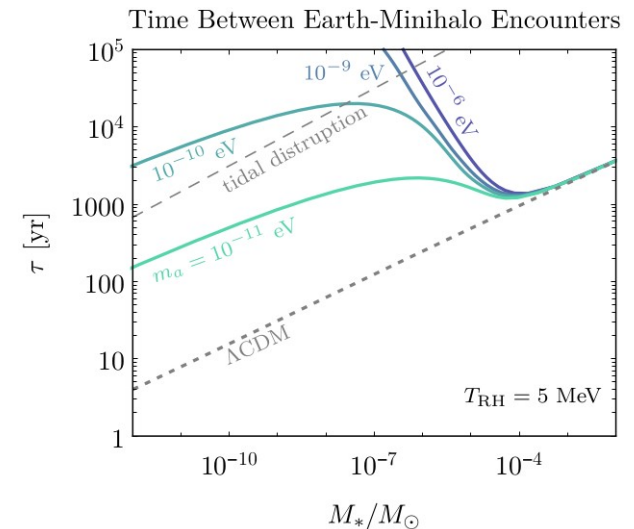
Direct Detection Implications

- Clumps can be disrupted by mergers, clump-clump encounters, encounters with the galactic disk, stars, etc.
- Structures collapsing before $z=250$ probably survive.

- Small couplings already difficult: $\propto 1/f_a$
- Time between earth-clump encounters can be larger than experimental lifetimes

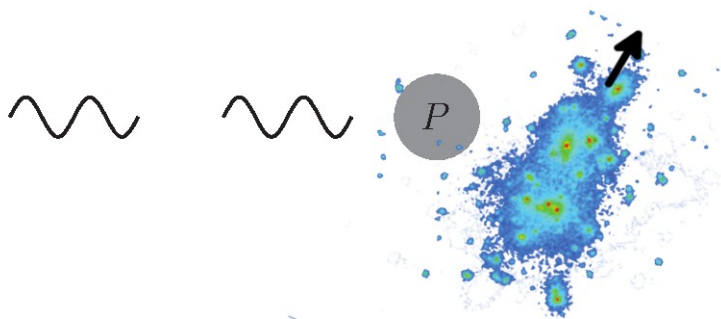
$$\tau = \frac{1}{n\sigma v}, \quad \text{where } n \sim \rho_0/M_*, \quad \sigma = \pi R_*^2$$

- Can create DM streams from clump disruption
- All fertile ground for N-body simulations.

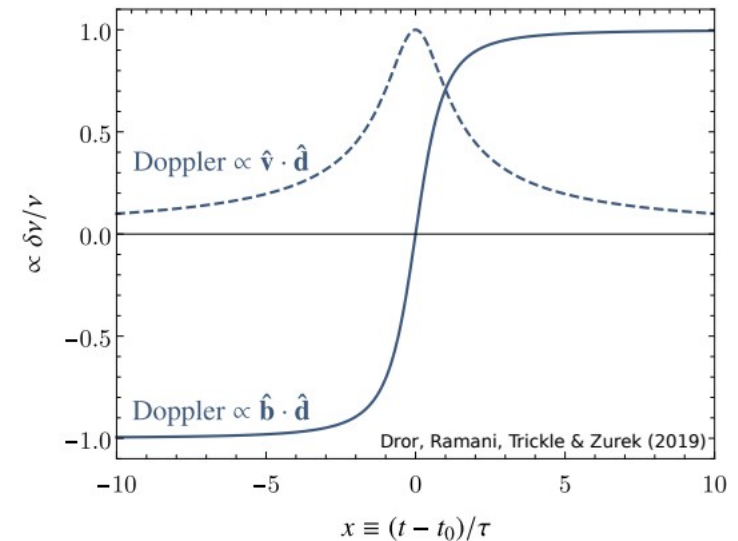


Astrophysical Probes: Pulsar Timing Arrays

- Clump passes near line of sight to a pulsar.
- Gravitational interactions alters pulse arrival time

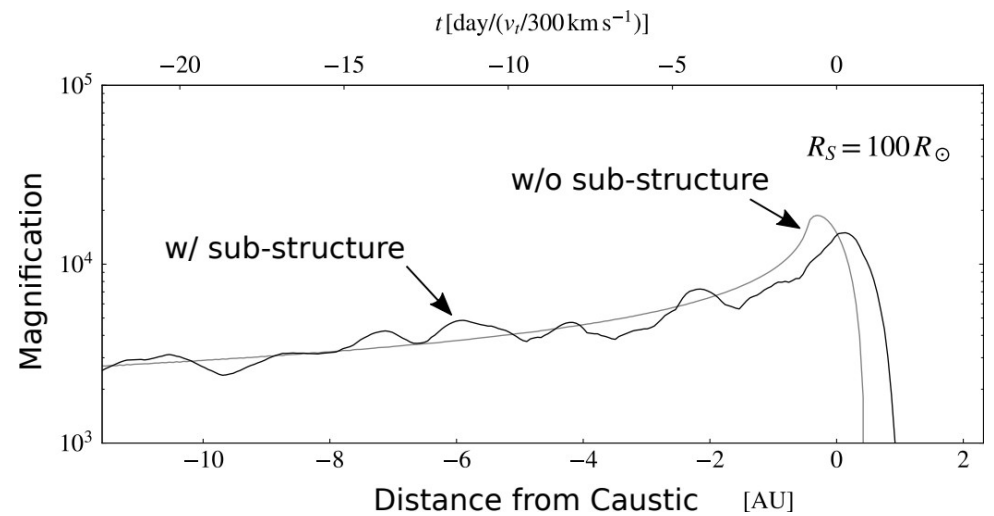
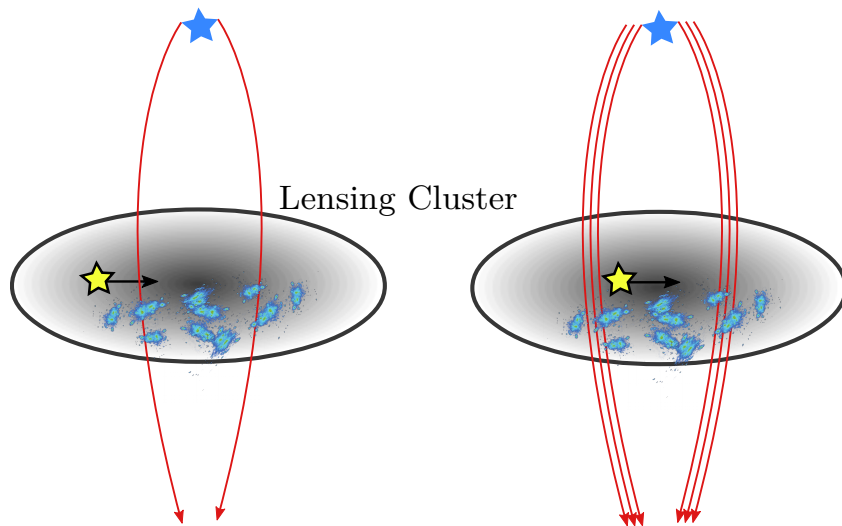


$$\frac{\delta\nu}{\nu} \sim \frac{GM}{vr_{min}}$$



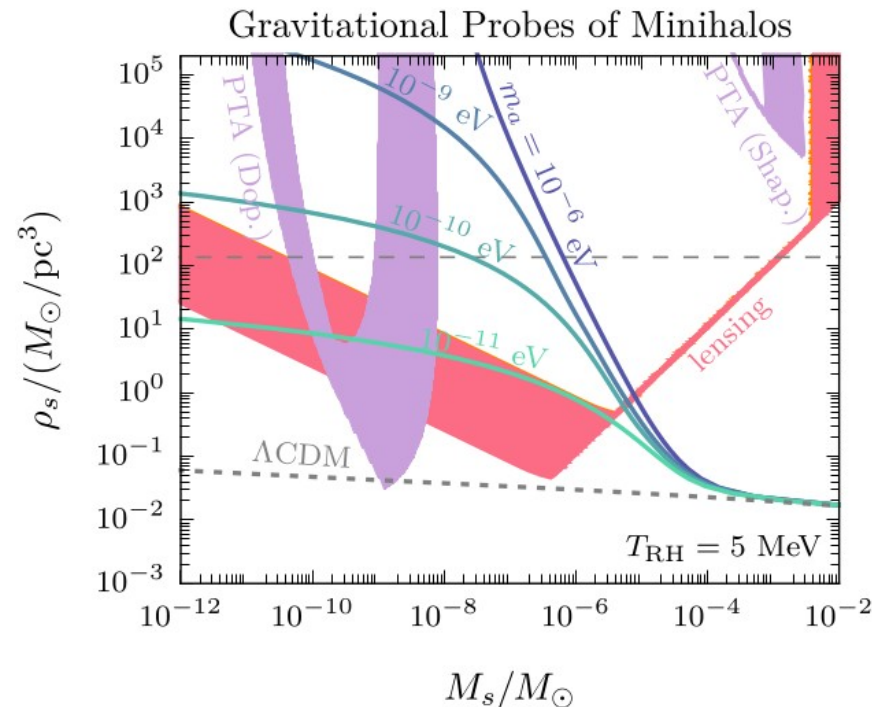
Astrophysical Probes: Caustic Microlensing

- Extragalactic star is lensed by a cluster, and then microlensed by an intra-lens star
- First known example in 2017 using Hubble (by accident).
- Tiny density fluctuations due to clumps amplified.
- Signature is 'noise' on the microlensing lightcurve.



Astrophysical Probes: Caustic Microlensing

- Assumes PTA of 200 pulsars observed for 20 years with SKA (but ignores important sources of noise!)
- Bands vary fraction of DM in clumps from 0.3 to 1.
- PTAs and Lensing may be sensitive to EMD-motivated very small-scale substructure.



Conclusions

- Non-thermal DM production sensitive to early universe cosmology.
- Well-motivated targets can be easier/harder for direct detection than standard scenario.
- A number of other possibilities lead to enhanced substructure from early universe dynamics.
 - Substructure is a window into the pre-BBN universe.

